

Supplementary Materials for
A single-phase epitaxially grown ferroelectric perovskite nitride

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Figs. S1 to S30
References

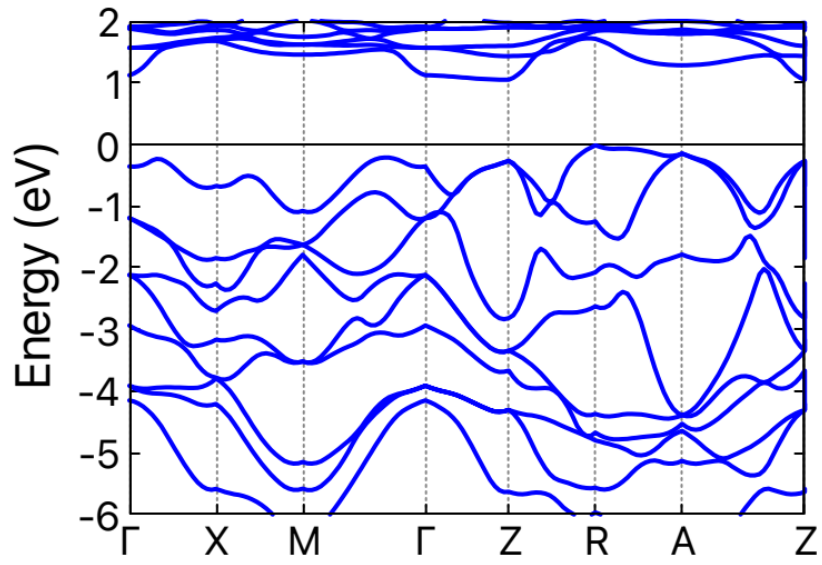


fig. S1. Band structures, as calculated with the HSE06 functional for CeTaN₃. The indirect band gaps are indicated. The zero of energy is set at the valence-band maximum.

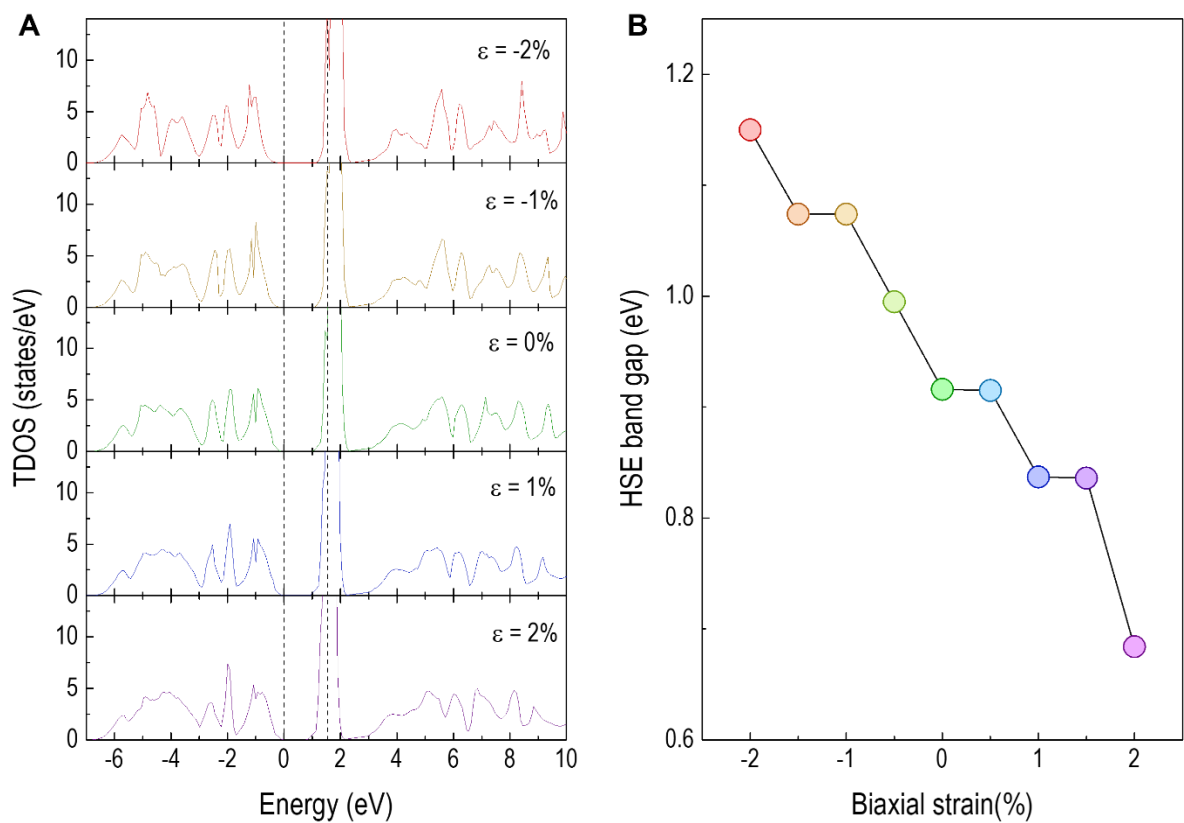


fig. S2. The band gap evolution of CeTaN₃ as a function of biaxial strain. (A) Total DOS of CeTaN₃ with biaxial strain ranging from -2% to 2% . Increasing epitaxial strain reduces the band gap, which is mainly caused by the downward movement of conduction band. (B) HSE bandgap evolution as a function of biaxial strain.

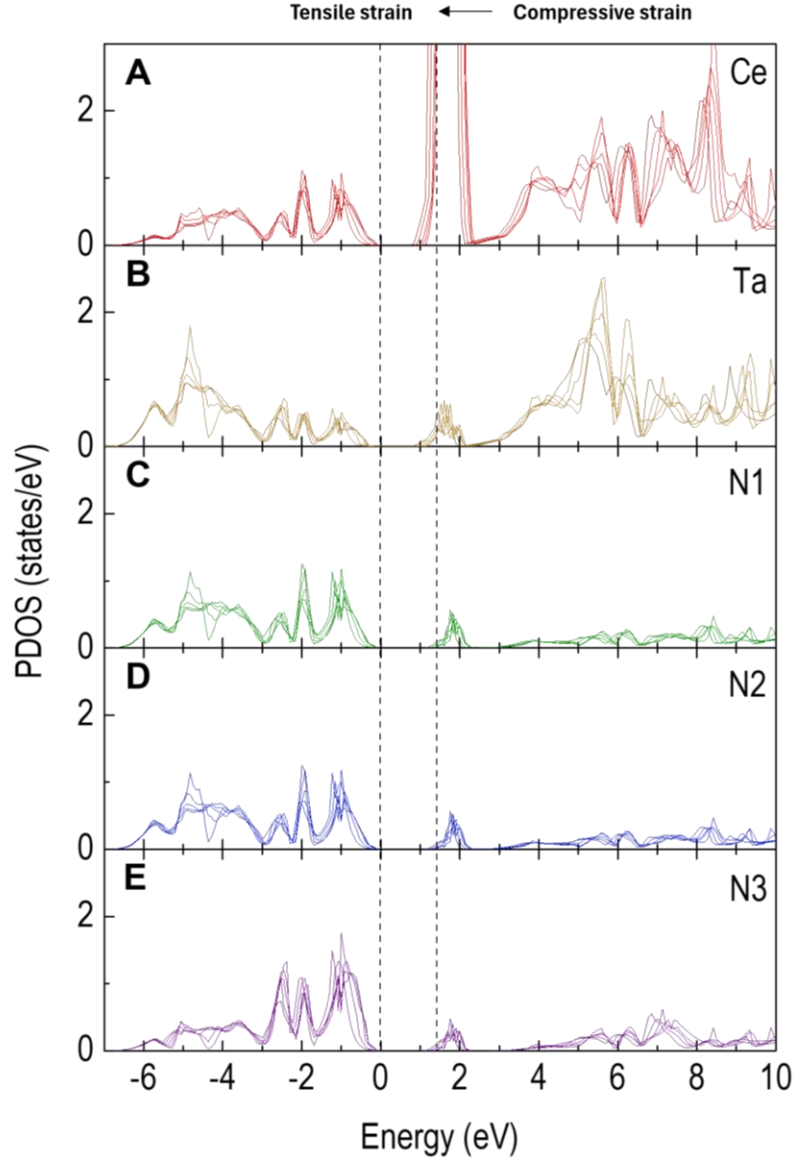


fig. S3. The comparison of PDOS of all elements under different strain states. PDOS of (A) Ce, (B) Ta, (C) N1, (D) N2, and (E) N3. Biaxial strain produces an additional energy shift to all atoms, leading to a smaller band gap under tensile strain and a larger band gap under compressive strain. Both tensile and compressive strain causes conduction band downward, while only compressive strain causes the valence band upward.

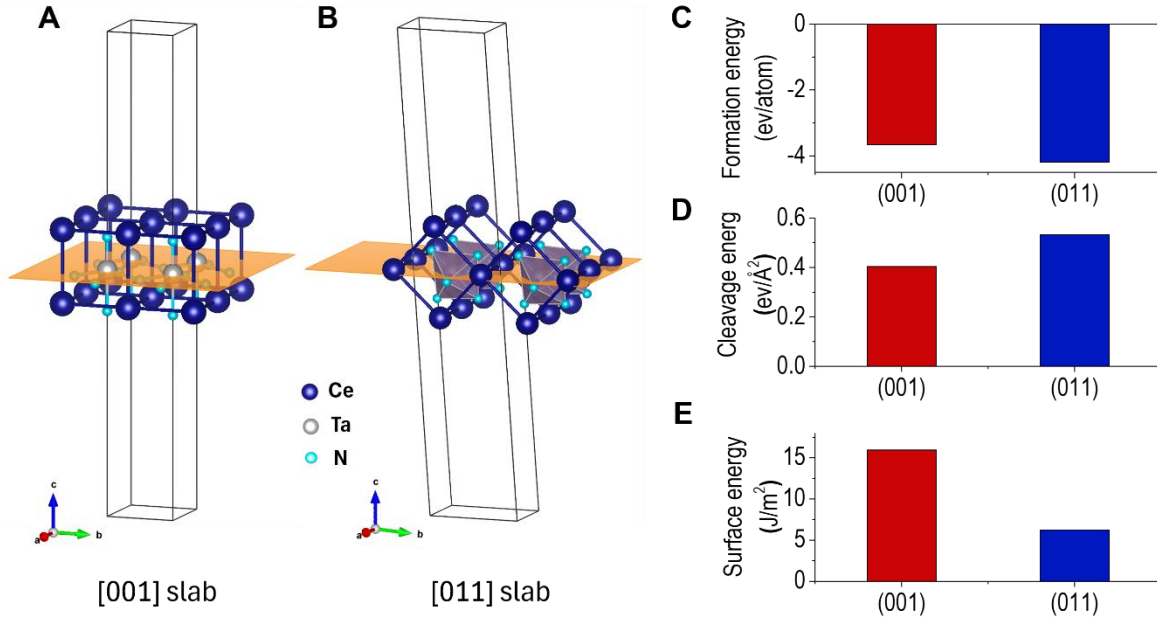


fig. S4. Comparison of formation energy, cleavage energy, and surface energy for (001)-, and (011)-oriented CeTaN₃. We can calculate the formation energy E_f by $(E_{slab} - n(E_{Ce} + E_{Ta} + 3 * E_N)) / N_{atom}$, the cleavage energy E_{cle} by $(E_{top}^{unrel} + E_{bottom}^{unrel} - E_{bulk}) / 2S$, and the surface energy E_{surf} by $(E_{slab} - nE_{bulk}) / A$, where E_{slab} is the total energy of the relaxed CeTaN₃, E_{bulk} is the total energy of CeTaN₃ bulk with the same crystal structure. E_{Ce} , E_{Ta} , and E_N are total energy of single atom of Ce, Ta, and N, respectively. E_{top}^{unrel} is the unrelaxed energy of the top layer after cleavage, E_{bottom}^{unrel} is the unrelaxed energy of the bottom layer after cleavage. All energies can be calculated using first-principles calculations based on density functional theory.

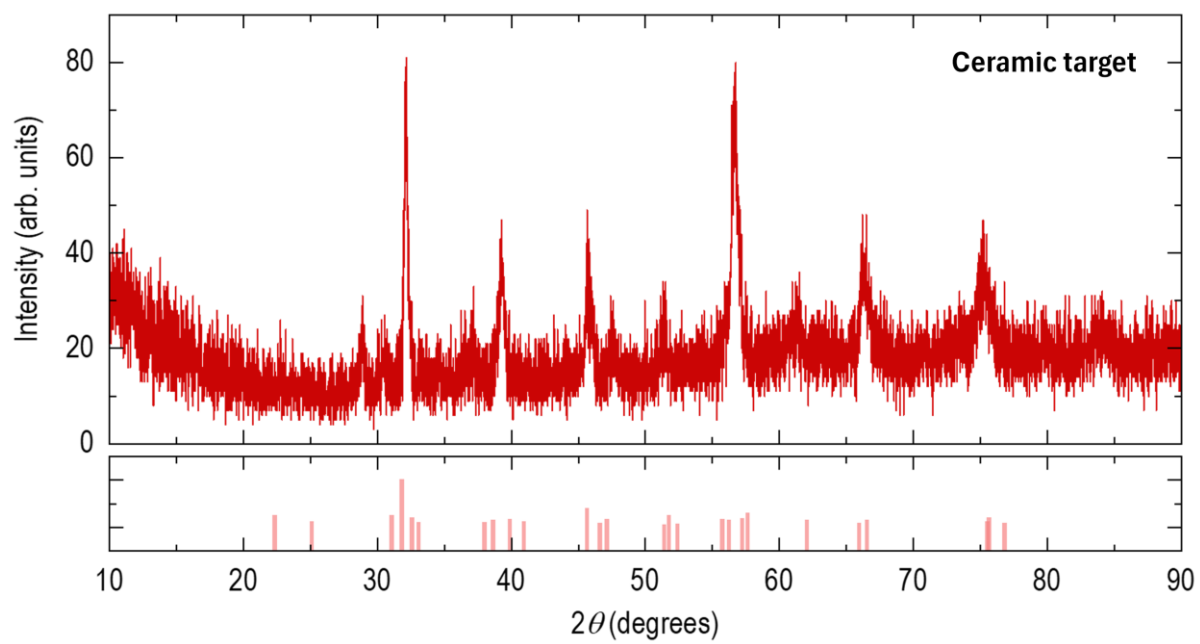


fig. S5. Powder X-ray diffraction θ - 2θ of CeTaN_3 ceramic target.

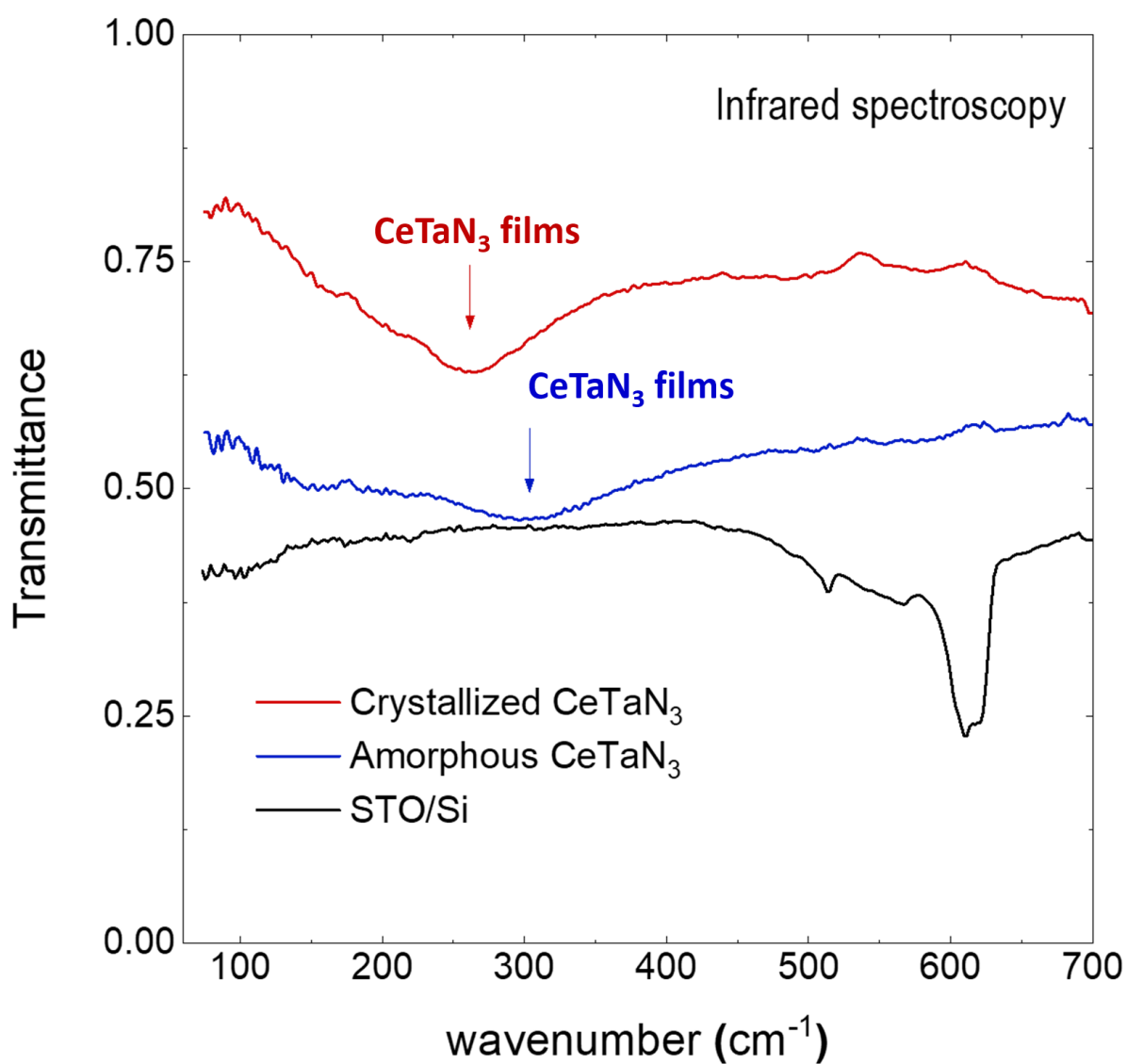


fig. S6. Infrared spectroscopy of crystalline CeTaN₃ with reference data from an amorphous CeTaN₃ and a SrTiO₃/Si substrate. Arrow lines mark the presence of CeTaN₃, in sharp difference with pure SrTiO₃/Si substrates.

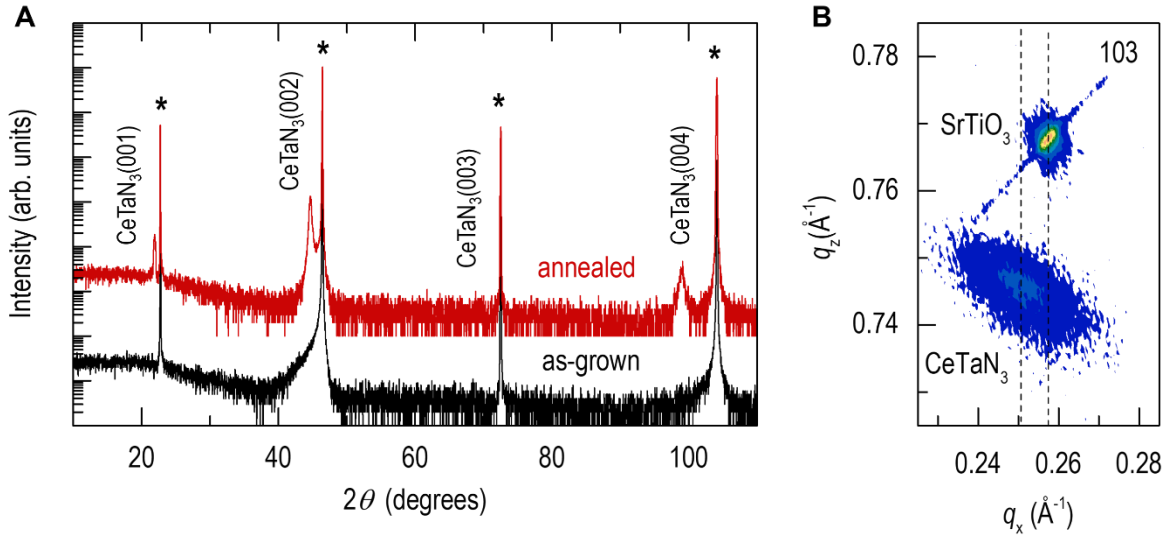


fig. S7. XRD measurements on amorphous and crystallized CeTaN_3 thin films on (001)-oriented SrTiO_3 substrates using rapid thermal process (RTP). (A) XRD θ - 2θ scans of amorphous CeTaN_3 and crystalline CeTaN_3 before and after RTP, respectively. RTP was performed at the temperature of 800°C in ammonia (NH_3) conditions with 10,000 Pa for a 1 hour with heating rate of $200^\circ\text{C}/\text{min}$. (B) RSM of CeTaN_3 around 103 peaks of SrTiO_3 , yielding the in-plane and out-of-plane lattice constants of CTN films are 3.99 and 4.02 Å, respectively.

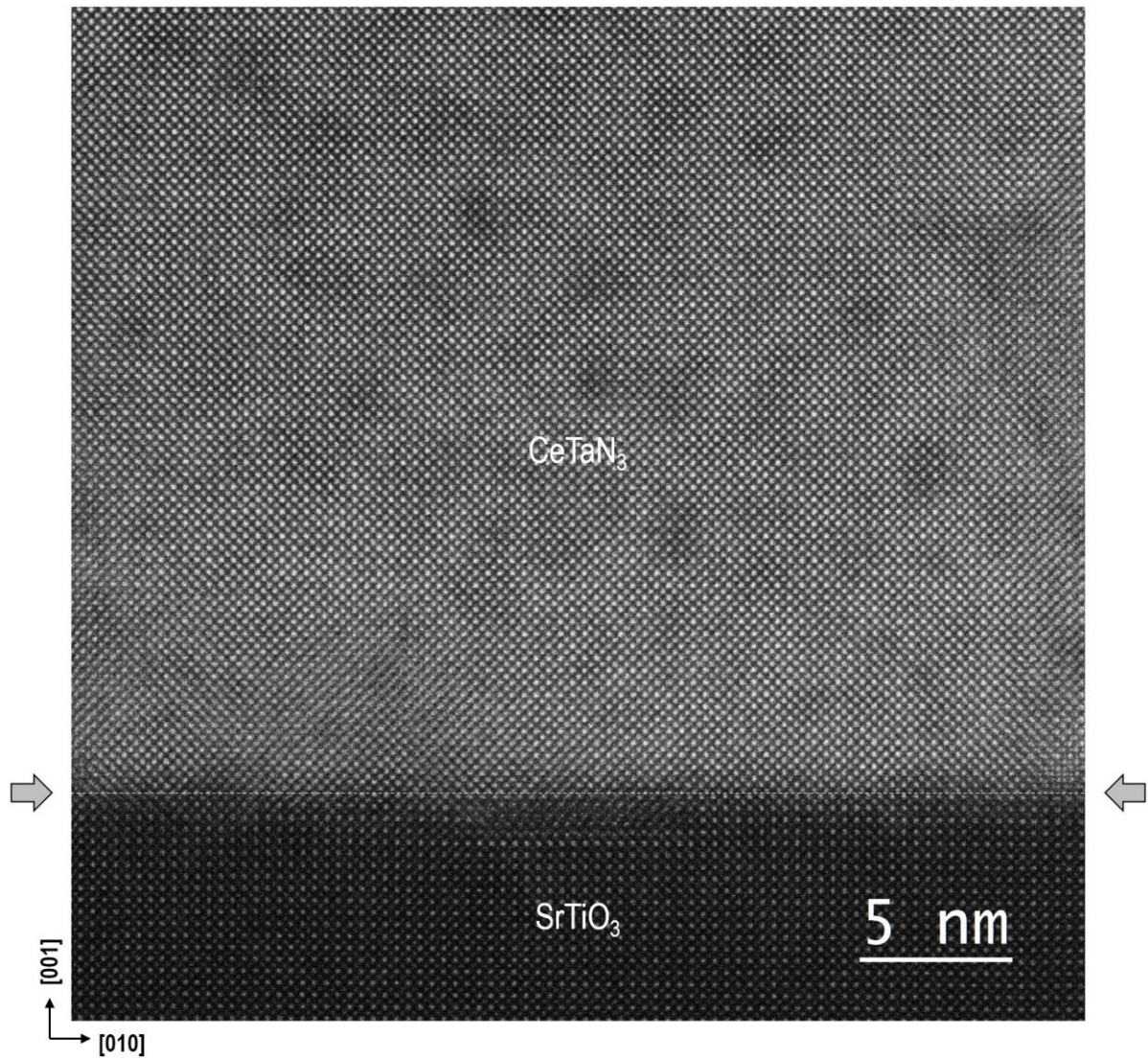


fig. S8. STEM-HAADF image of a CeTaN₃ thin film grown on SrTiO₃ substrates. Dashed line and arrows indicate the interfaces between CeTaN₃ and SrTiO₃. The sharp interfaces and well-ordered atoms indicate that the CeTaN₃ thin films have high crystallinity.

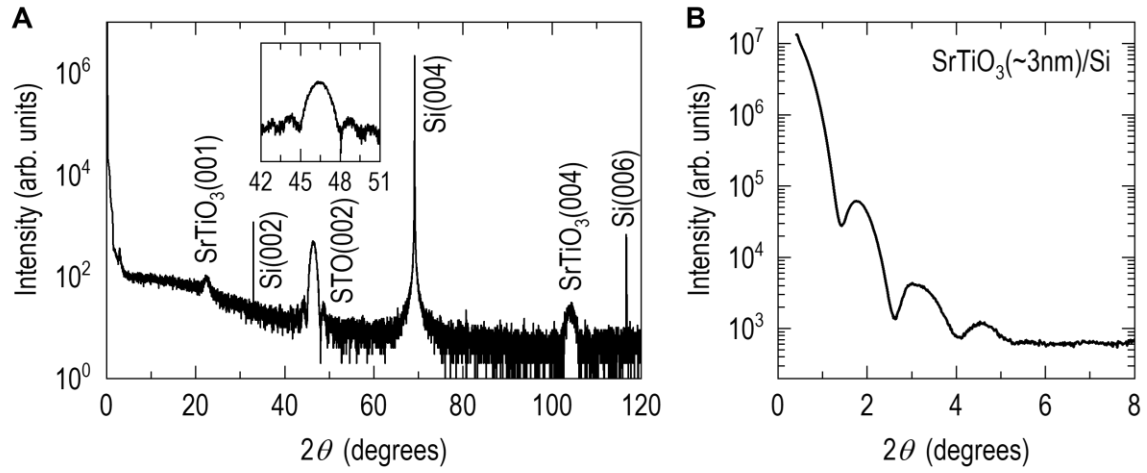


fig. S9. Structural characterization of 3 nm-thick STO capped Si substrate. (A) XRD θ - 2θ scan exhibits the epitaxial growth of STO on Si without any impurity peaks. The inset shows the detail of the STO 002 peak, revealing the Kiessig fringes. (B) X-ray reflection (XRR) shows Fourier oscillations, confirming the flat surface and about 3nm thickness of the STO capping layer.

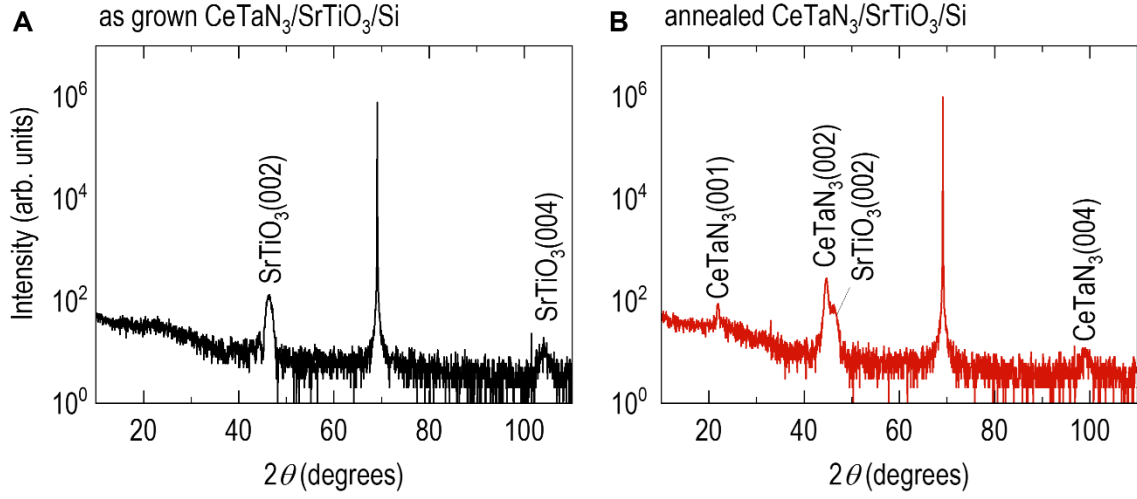


fig. S10. XRD θ - 2θ of CeTaN_3 crystallization on STO/Si substrate using RTP. (A) Before RTP, XRD result only shows the STO peaks without any peaks from CeTaN_3 films, suggesting that the CeTaN_3 is amorphous. (B) We find the $00l$ reflections of CeTaN_3 films, indicating that the CeTaN_3 films are crystallized after RTP.

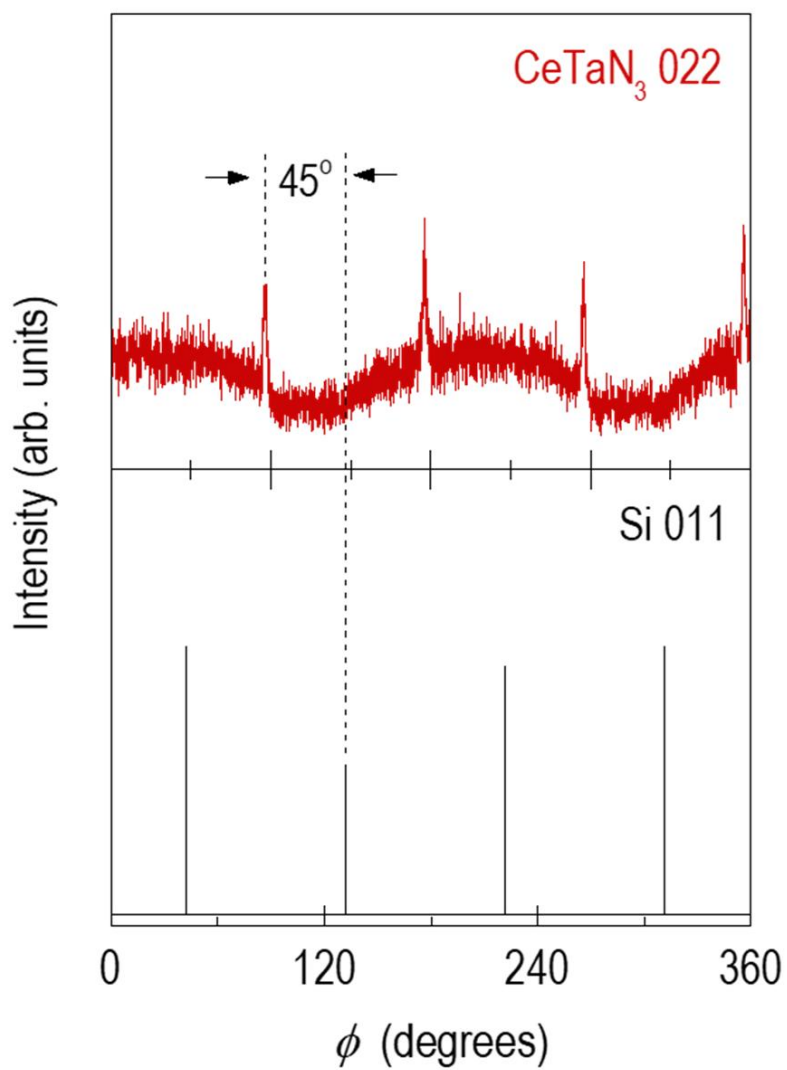


fig. S11. Phi-scans for CeTaN_3 022 and Si 011 reflections. Two sets of curves shift by 45 degrees.

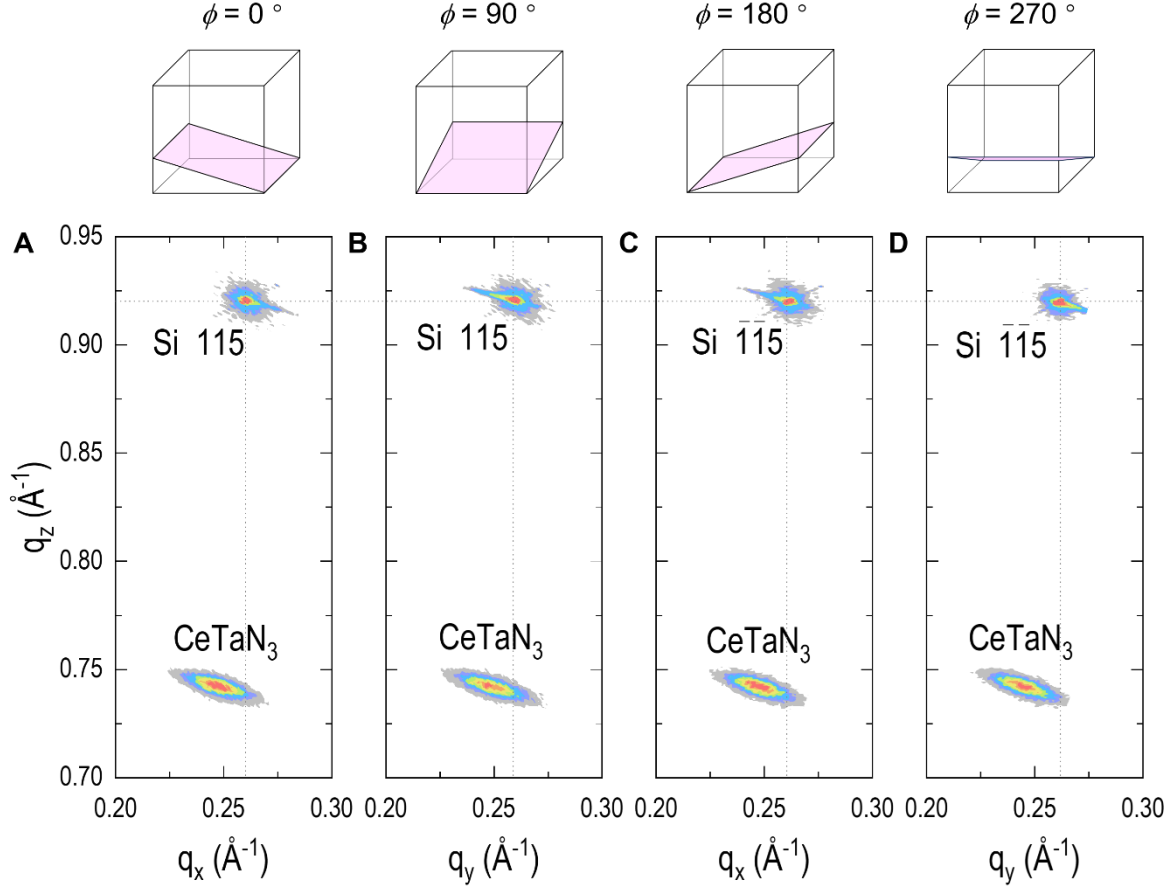


fig. S12. Tetragonal structure of CeTaN₃ thin films confirmed by RSM. The diffraction peaks, associated with the indices 103, 013, $\bar{1}03$, and $0\bar{1}3$ reflections of CeTaN₃, reveal consistent positioning of q_x and q_y . The phi-dependent RSM results suggest that the CeTaN₃ thin film exhibits a tetragonal structure. The average in-plane lattice parameter is calculated to be $a = b = 4.03 \text{ \AA}$, while the averaged out-of-plane lattice parameter is $c = 4.09 \text{ \AA}$.

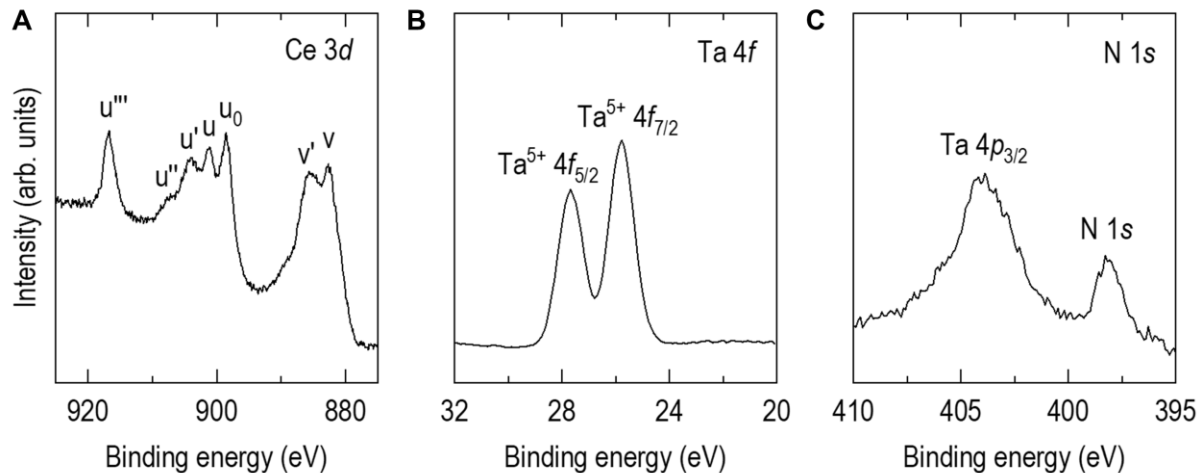


fig. S13. X-ray photoelectron spectroscopy of CeTaN₃ thin films. Ce 3d, Ta 4d, and N 1s, respectively. (A) Ce represents two oxidation states of Ce³⁺ and Ce⁴⁺. The peaks of u' (901.3 eV), u₀ (898.5 eV), and v' (885.5 eV) correspond to the positions of Ce³⁺ determined from Ce₂O₃, while peaks of u''' (916.8 eV), u'' (907.2 eV), u (901.1 eV) and v (882.7 eV) match Ce⁴⁺ determined from CeO₂. The presence of two oxidization states may attribute to the inevitable nitrogen vacancies in the surface layers of CeTaN₃ thin films. (B) The result of Ta spectroscopy reveals a single oxidation state of Ta⁴⁺. Ta⁵⁺ 4f_{5/2} is located at 27.7 eV, and Ta⁵⁺ 4f_{7/2} is at 25.8 eV. (C) N shows a distinct single peak at 398.1 eV, nearby where a peak corresponding to Ta 4p_{3/2} was observed at 403.9 eV.

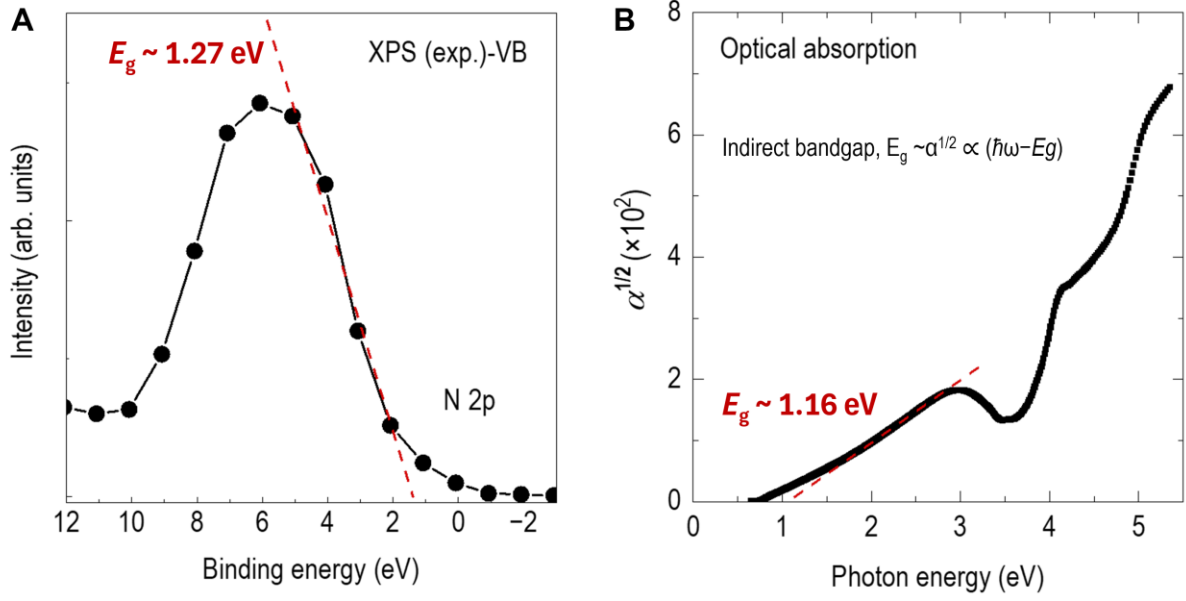


fig. S14. Determination of band gap of CeTaN₃. (A) XPS valence band (VB) spectra for CeTaN₃ thin films. The band gap estimated from VB spectra is approximately 1.27 eV. (B) Optical absorption spectra of a CeTaN₃ film on double-polished SrTiO₃ substrate. α represents the absorption coefficient. We plotted the $\alpha^{1/2}$ against the photon energy ($h\omega$). The linear relationship between $\alpha^{1/2}$ and $h\omega$ demonstrates that the CeTaN₃ is an indirect band gap semiconductor, consistent with theoretical calculations. The optical band gap determined from absorption spectra is approximately 1.16 eV.

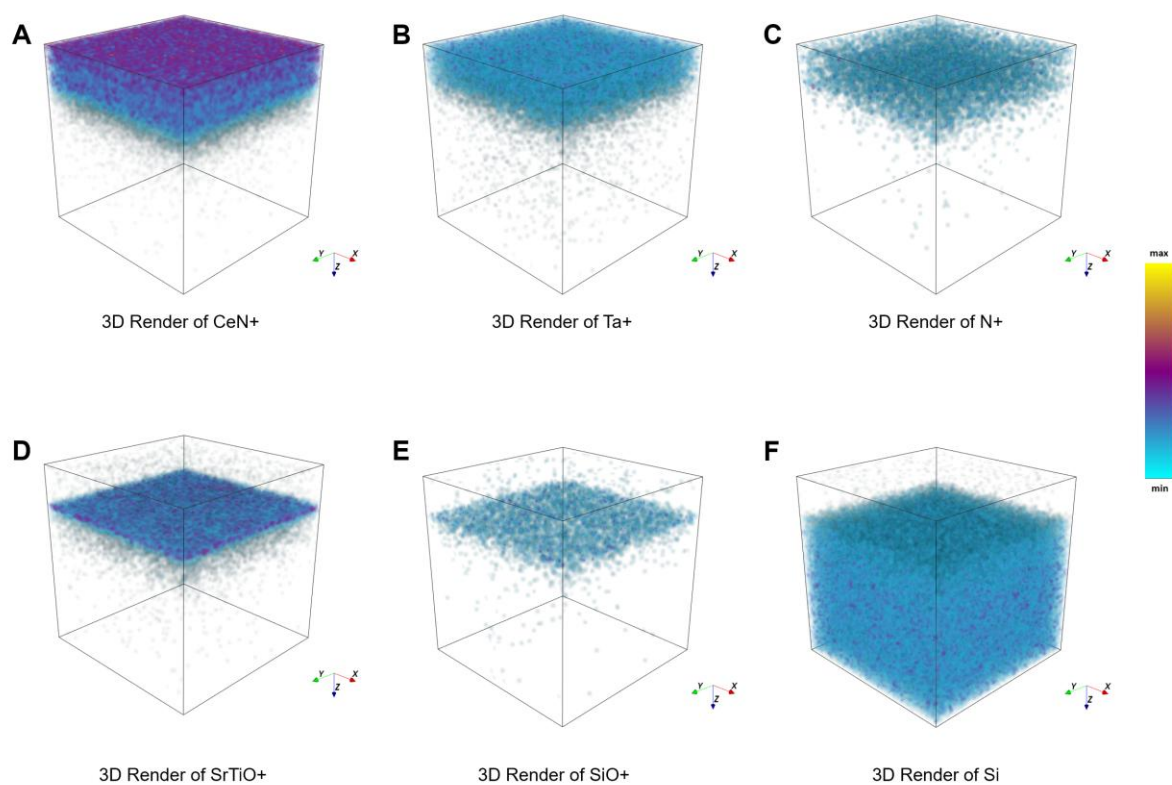


fig. S15. ToF-SIMS analysis of $\text{CeTaN}_3/\text{SrTiO}_3/a\text{-SiO}_2/\text{Si}$. (A–F) 3D renderings of CeN^+ , Ta^+ , N^+ , SrTiO^+ , SiO^+ , and Si signals, respectively. The SIMS results confirm the uniformity of all deposited layers and provide insight into the chemical distribution from the top to the bottom of the structure.

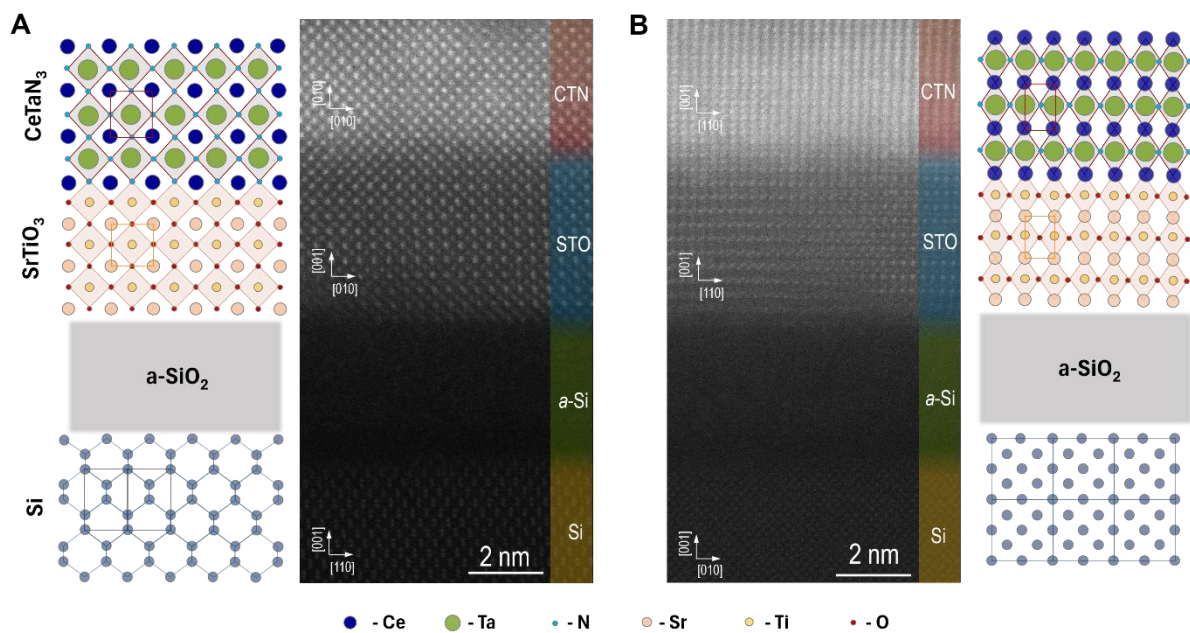


fig. S16. High resolution STEM images for a CeTaN₃/SrTiO₃/a-SiO₂/Si sample. (A) and (B) illustrate the STEM images viewed from Si [110] and Si [010] orientations, respectively.

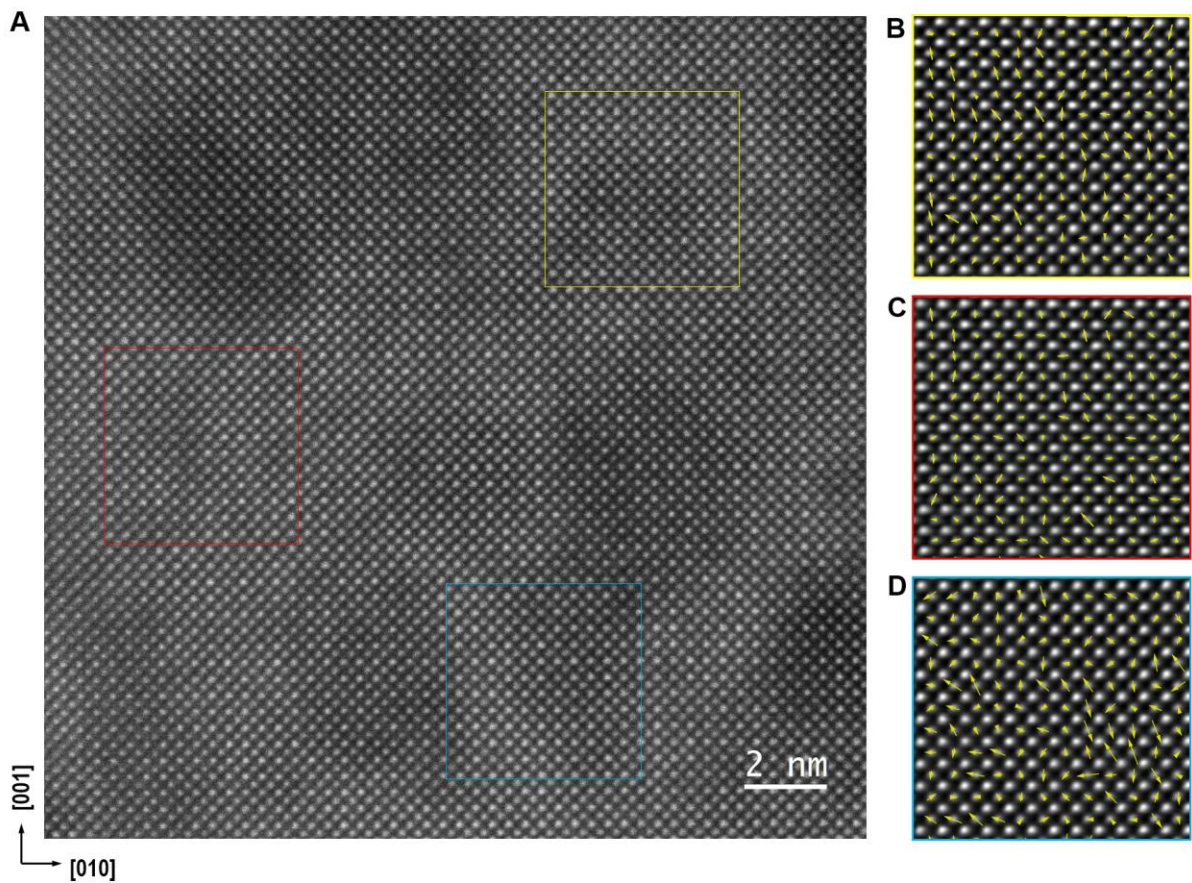


fig. S17. Polarization analysis of CeTaN₃. (A) A representative of zoom-in HAADF image of CeTaN₃. (B to D) The Ce atoms displacements were deduced from three representative TEM images. The arrows representing the motion of Ce atoms results to partial polarization, with most of these arrows pointing upward and some pointing to the left.

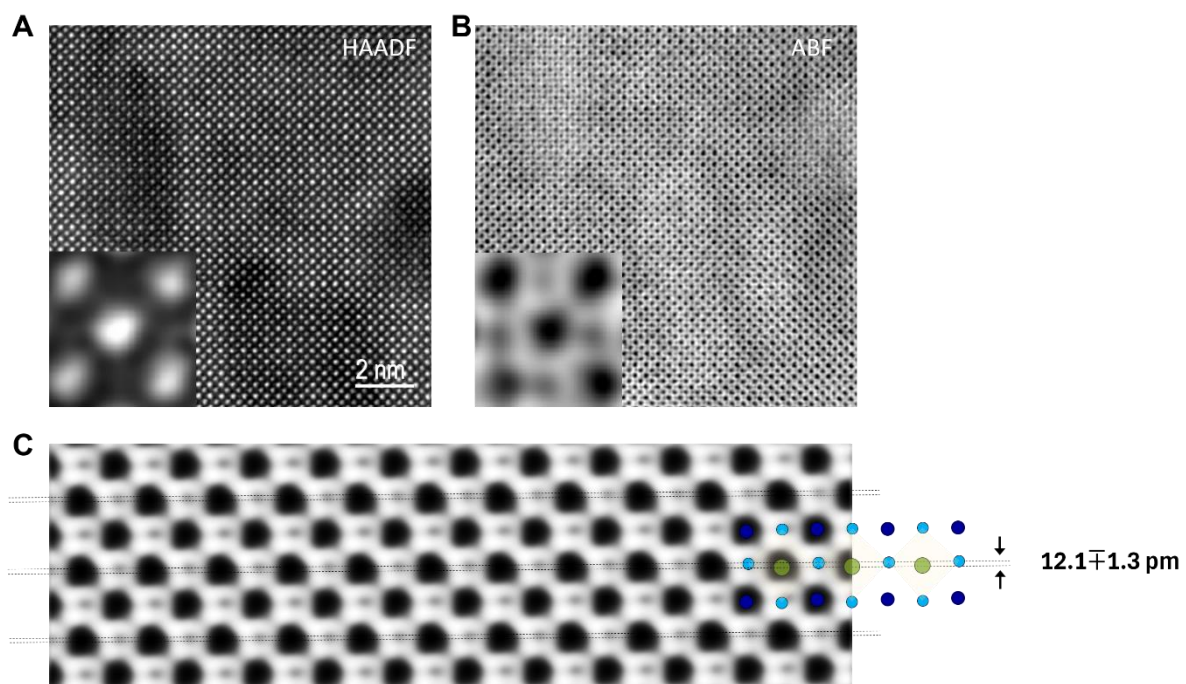


fig. S18. STEM results of CeTaN_3 thin films. To investigate atomic displacements intuitively related with ferroelectricity, (A) HAADF (high angle annular dark field) and (B) ABF (annular bright field) STEM images were recorded. (C) A ABF-STEM image of CeTaN_3 , exhibiting schematic of non-centrosymmetric Ta atoms. The atomic displacement shifts from center position by 12.1 ± 1.3 pm.

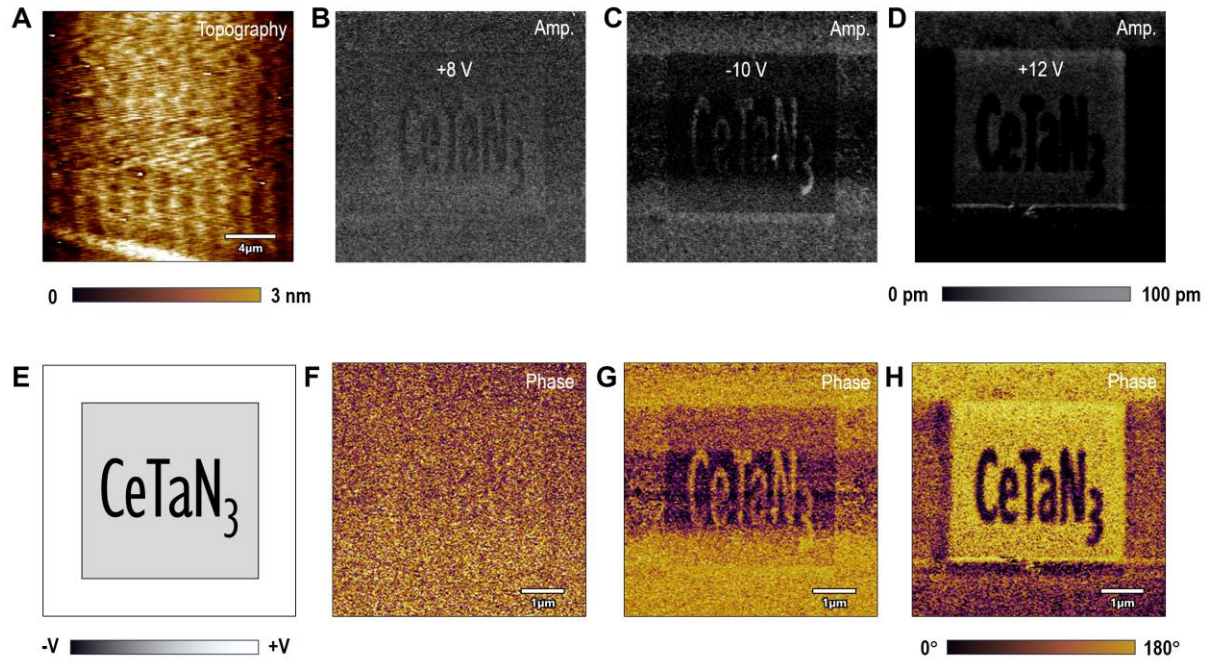


fig. S19. PFM results on a 15-nm-thick CeTaN_3 thin film grown on SrTiO_3/Si . (A) Topography of sample, yielding a r.m.s of ~ 2.7 nm. The “ CeTaN_3 ” letters were written with opposite tip voltages comparing to the switching voltages poling the square area. (B to D) PFM amplitude and (F to H) PFM phase images with gradually increasing tip voltages. (E) Schematic of electrical switching patterns.

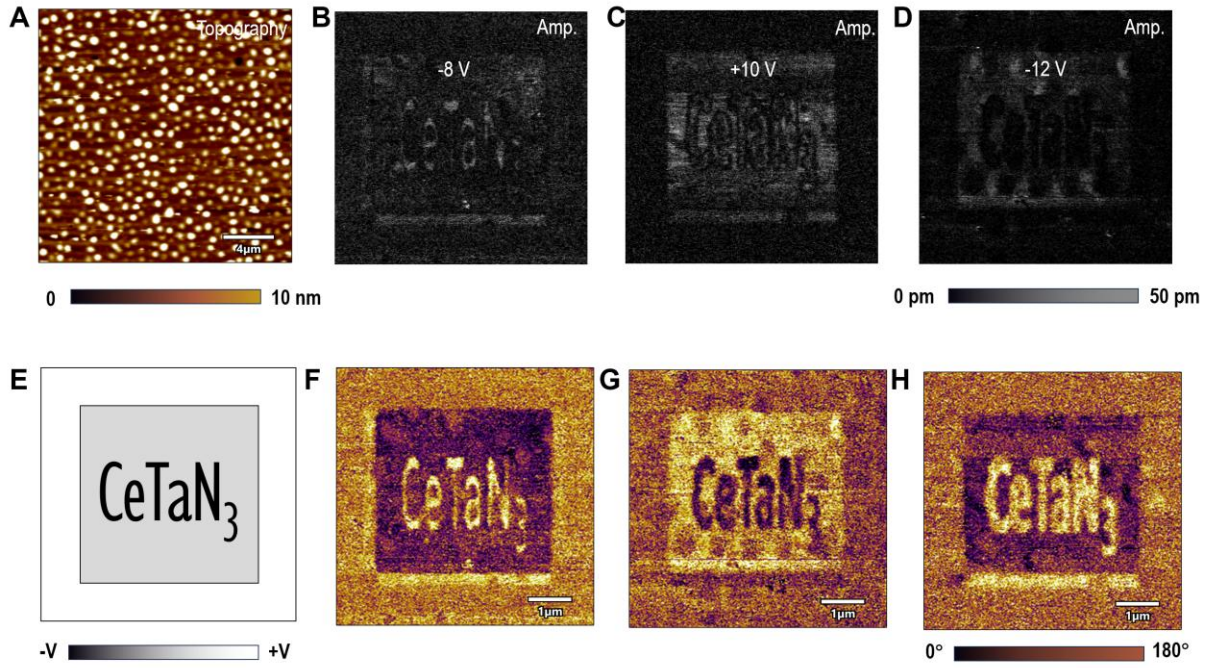


fig. S20. PFM results on an 80-nm-thick CeTaN_3 thin film grown on SrTiO_3/Si . (A) Topography of sample, yielding a r.m.s of ~ 8.3 nm. The “ CeTaN_3 ” characters were written with opposite tip voltages comparing to the switching voltages poling the square area. (B to D) PFM amplitude and (F to H) PFM phase images with gradually increasing tip voltages. (E) Schematic of electrical switching patterns.

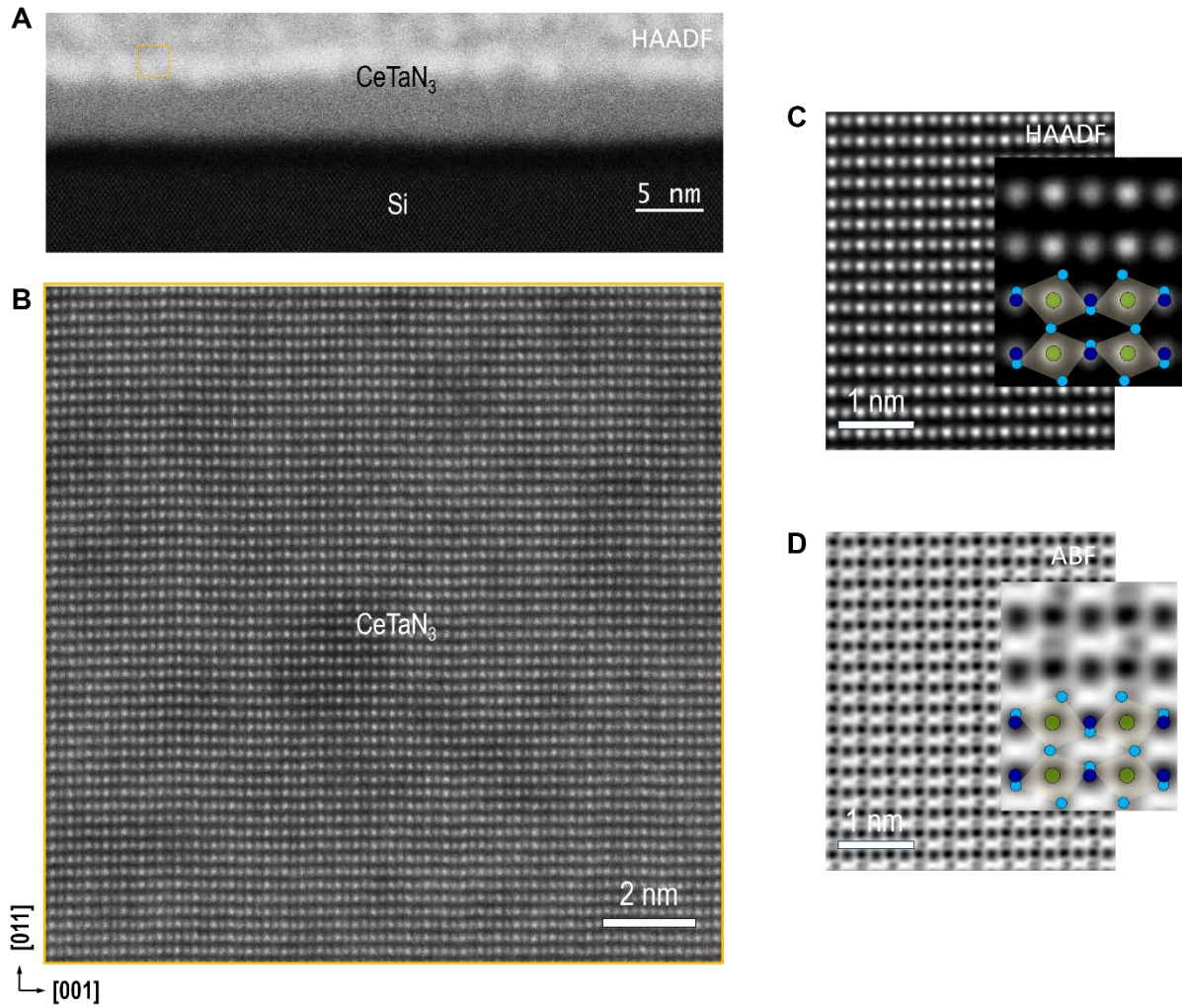


fig. S21. STEM analysis of CeTaN₃ thin films directly on Si. (A) Low magnification cross-sectional STEM-HAADF image at CeTaN₃/Si interfaces. Still, we cannot avoid the oxidation of Si at the interfaces at our growth conditions. The thickness of SiO₂ layer is approximately 1 nm. (B) High magnification STEM-HAADF image of CeTaN₃ thin films. The out-of-plane direction is along the [011] orientation. (C) STEM-HAADF and (D) STEM-ABF images of CeTaN₃, from which we could identify the tilted nitrogen octahedral.

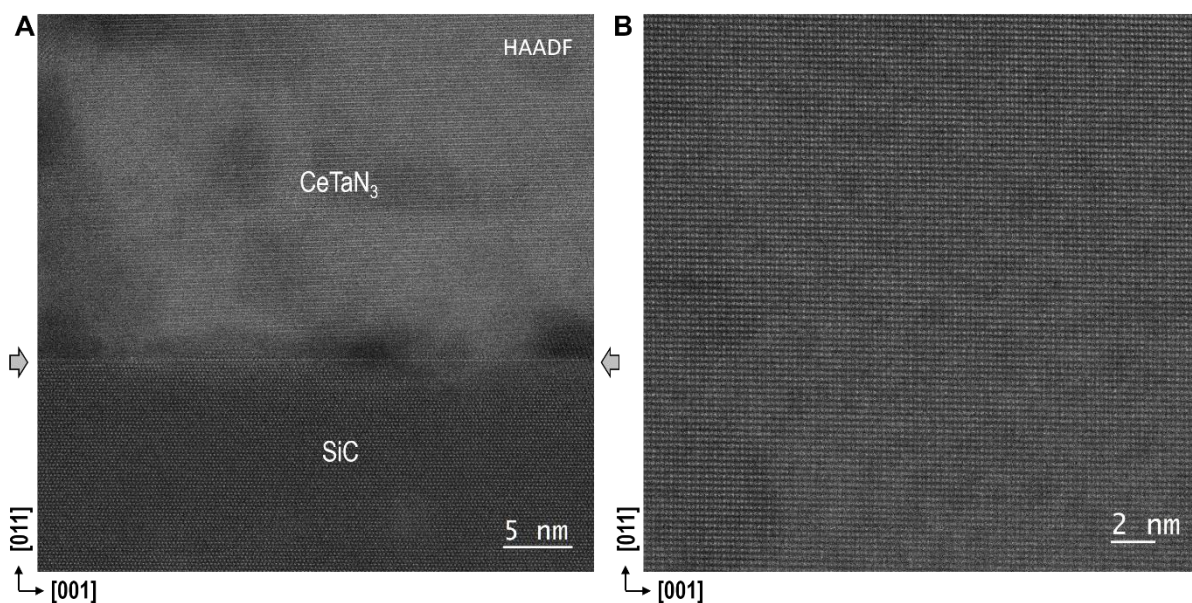


fig. S22. STEM analysis of CeTaN₃ thin films directly on SiC. (A) STEM-HAADF image at the CeTaN₃/SiC interface region. Arrows and dashed line indicate the position of interface. (B) High magnification STEM-HAADF image of CeTaN₃ thin films.

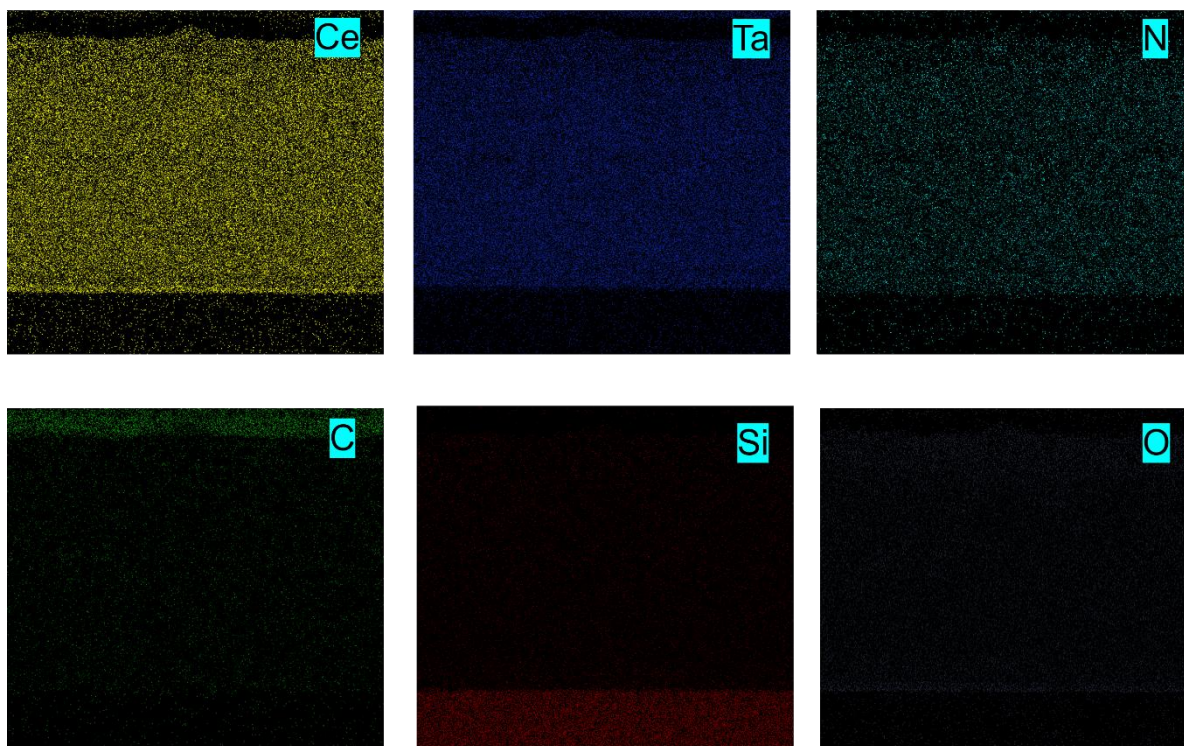


fig. S23. Chemical distribution in CeTaN₃ thin films. (A to F) EDX images of Ce, Ta, N, C, Si, and O elements. These results indicate that the CeTaN₃ films exhibit a high-crystallinity and chemical homogeneity without apparent contamination of oxygen.

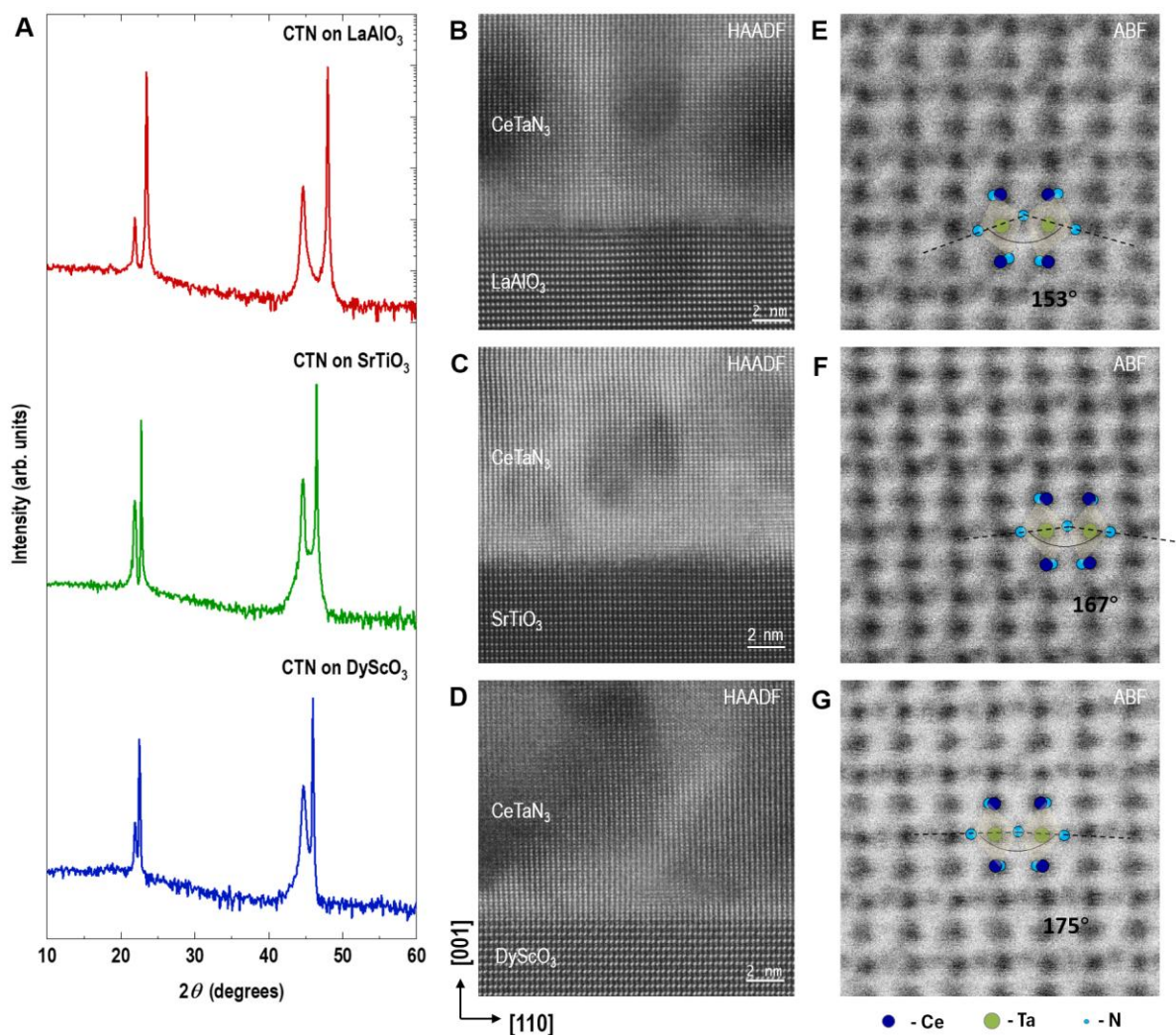


fig. S24. Structural characterization of CeTaN_3 thin films on different oxide substrates. (A) XRD θ - 2θ scans of CeTaN_3 thin films grown on LaAlO_3 , SrTiO_3 , and DyScO_3 substrates, confirming their epitaxial growth. (B–D) HAADF-STEM images of CeTaN_3 thin films on LaAlO_3 , SrTiO_3 , and DyScO_3 substrates, respectively, revealing atomically sharp interfaces. (E–G) High-magnification ABF-STEM images of CeTaN_3 thin films on LaAlO_3 , SrTiO_3 , and DyScO_3 substrates, respectively. Insets in each ABF image illustrate the schematic of octahedral tilting, with the corresponding tilt angles indicated.

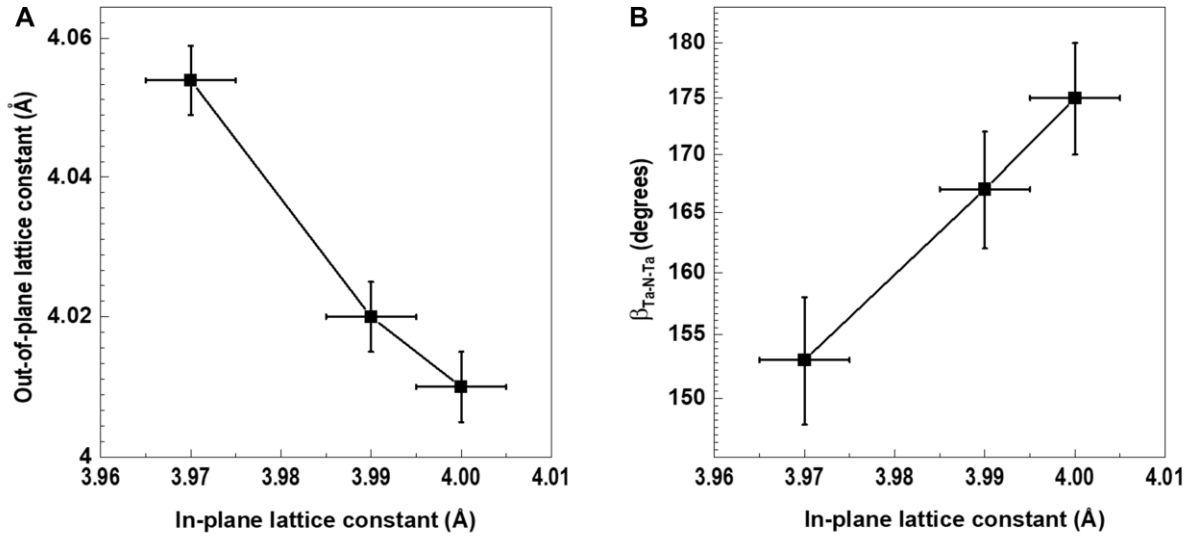


fig. S25. Structural parameter evolution with in-plane lattice constant. We determined the out-of-plane and in-plane lattice parameters of CeTa₂N₃ thin films using XRD measurements, including θ -2 θ scans and reciprocal space mappings (RSMs). **(A)** Out-of-plane lattice parameters and **(B)** $\beta_{\text{Ta-N-Ta}}$ angle as a function of in-plane lattice constants for three CeTa₂N₃ thin films. The results clearly show that the $\beta_{\text{Ta-N-Ta}}$ angle progressively decreases with increasing compressive strain, following the trend from DyScO₃ to SrTiO₃ and LaAlO₃ substrates.

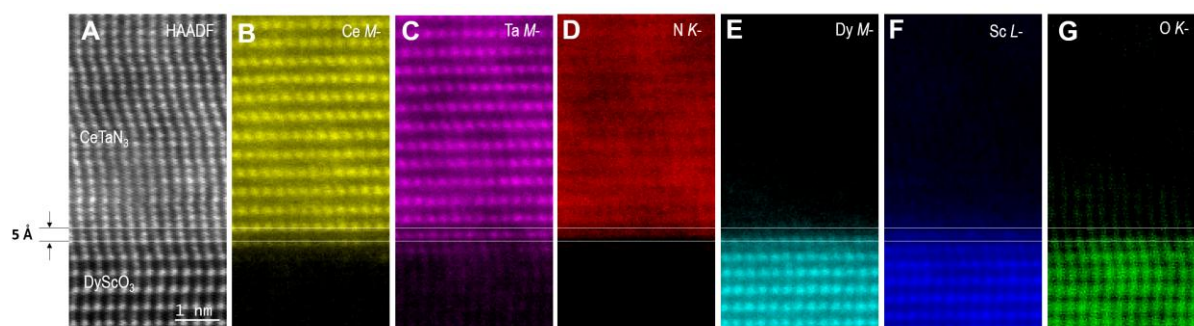


fig. S26. EELS mapping of CeTaN₃ thin films on DyScO₃ substrates. (A) HAADF-STEM image, where the tilt is attributed to thermal drift during STEM imaging. (B–G) Atomically resolved EELS maps corresponding to the Ce *M*-, Ta *M*-, N *K*-, Dy *M*-, Sc *L*-, and O *K*-edges. These results confirm the exceptionally high chemical homogeneity of CeTaN₃ thin films, with minimal chemical intermixing at the heterointerface, limited to a length scale of one unit cell. The interface termination is identified as TaN₂–DyO.

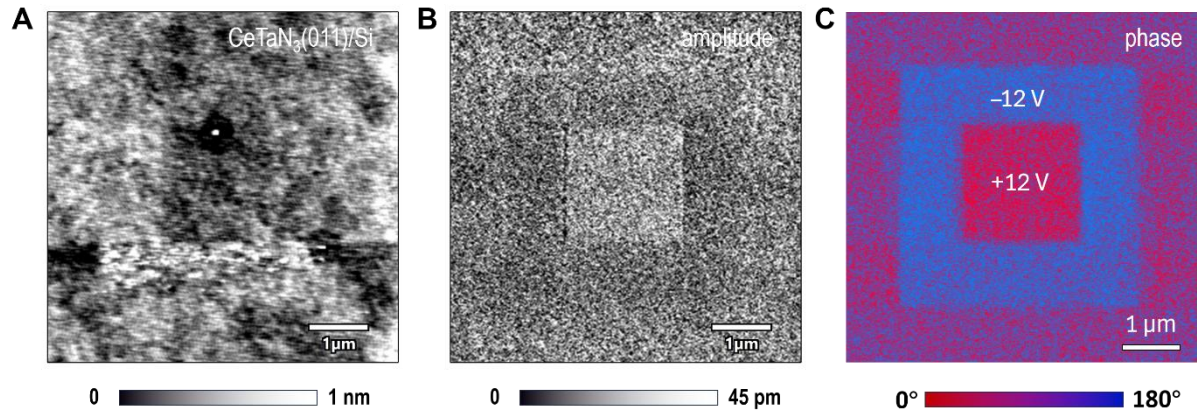


fig. S27. PFM results of CeTaN_3 thin films directly grown on Si. (A) Topography of CeTaN_3/Si . The r.m.s of sample is ~ 1 nm. (B) PFM amplitude and (C) phase images of CeTaN_3/Si . The center and border square area was upward and downward polarized by tip voltages, respectively. the square shapes show reversal phase and amplitude of piezoelectric responses.

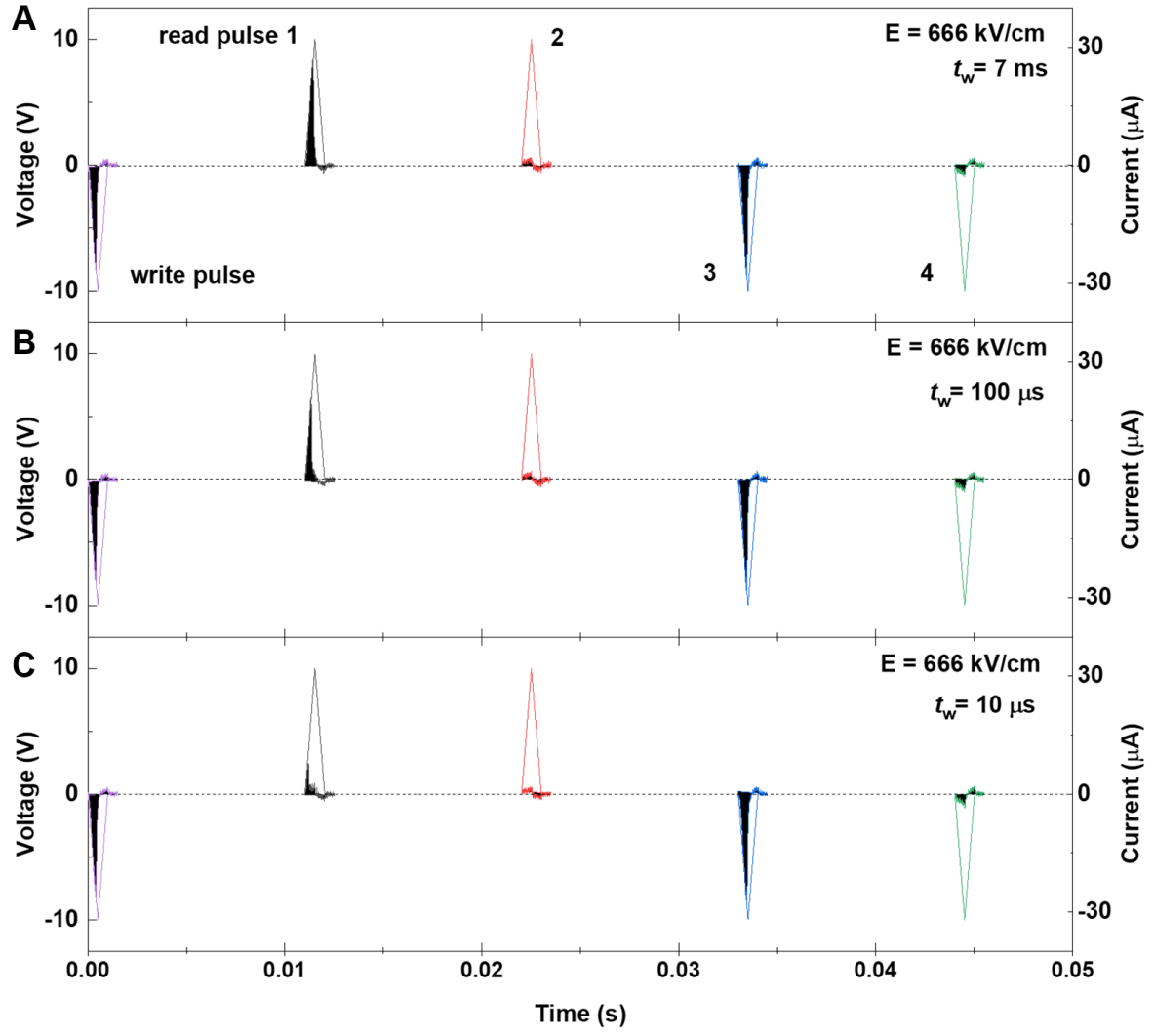


fig. S28. PUND measurements on a Pt/CeTaN₃/p-Si capacitor. (A–C) PUND measurements were conducted with a fixed poling electric field ($V_w = 10$ V, corresponding to 667 kV/cm) at different writing times ($t_w = 7$ ms, 100 μs, and 10 μs), respectively. A sharp difference is observed in the read pulse 1, revealing distinct switching currents.

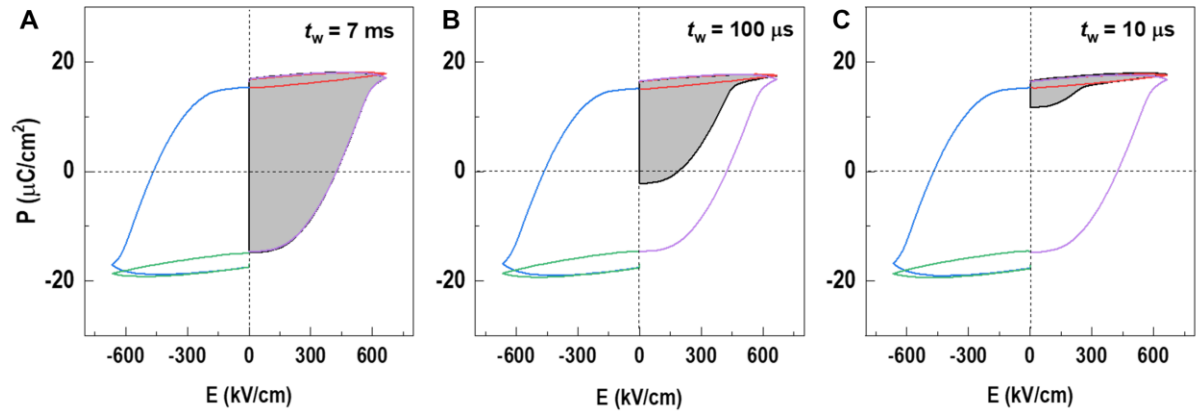


fig. S29. Ferroelectric hysteresis loops of Pt/CeTaN₃/p-Si capacitor under different writing times (t_w). These results are calculated from Fig. S27 A-C. The write and read curves correspond to the curves described in Fig. S27. The shadow areas represent the switched ferroelectric polarizations under different t_w .

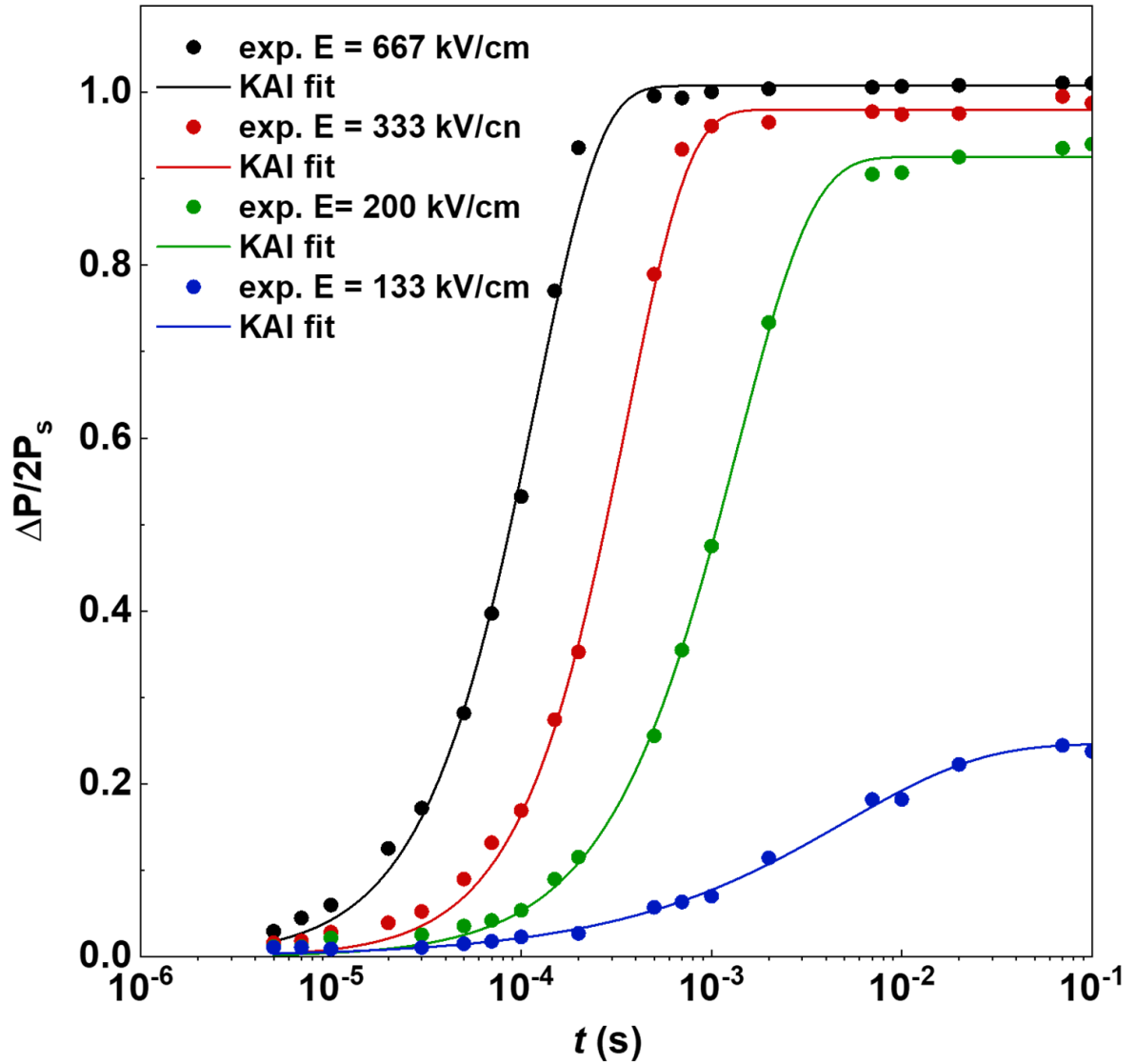


fig. S30. Ferroelectric switching tests on a Pt/CeTaN₃/p-Si capacitor. The switched polarization (ΔP) as a function of t for four different electric fields. ΔP were obtained by varying the pulse duration (t) while keeping the electric field constant. According to the Kolmogorov-Avrami-Ishibashi (KAI) model (58, 59), the switched polarization follows the relation $\Delta P = 2P_s\{1 - \exp[-t/t_{sw}]\}$ for a continuous one-dimensional medium. Fitting the data using the KAI model reveals that the switching time ranges from 1.18×10^{-4} s to 0.005 s, depending on the poling fields. Higher electric fields result in faster ferroelectric domain switching.

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